

Towards Shorter Wavelength EUV and Soft X-ray Sources

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Outline

- Background (Sn and Xe)
- Results at high Z (W) and implications for other elements
- Recent work on Gd and Tb sources for 6.7 nm
- Other potential sources







Sn is the brightest source at 13.5 nm



Because of opacity, the UTA narrows and increases in intensity with decreasing tin concentration (up to $^{2\%}$).

As concentration decreases out-of-band emission decreases.

Max inband conversion efficiency ~ 2.3% per 2π sr for 100% Sn at $\phi \approx 1.6 \times 10^{11}$ Wcm⁻² ~ 2.9% per 2π sr for 5% tin at $\phi \approx 2 \times 10^{11}$ Wcm⁻²











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Origin of the Sn UTA

Levels of 4p⁶4dⁿ & (4p⁵4dⁿ⁺¹+4dⁿ⁻¹4f + 4dⁿ⁻¹5p) Calculated with Cowan suite of codes



The Sn UTA is due to $4p^{6}4d^{n} \rightarrow (4p^{5}4d^{n+1}+4d^{n-1}4f + 4d^{n-1}5p) (0 \le n \le 9)$ transitions. CI (configuration interaction) causes transitions in successive ion stages to partially overlap in energy.

(O'Sullivan and Faulkner Opt. Eng. 33, 3978-3983 (1994), Koike et al J. Elec. Spec. Relat. Phen. 144, 1227 (2005)

CI effects also cause a spectral narrowing, so the individual arrays have widths of ~ 5eV.





Sn⁵⁺–Sn¹³⁺ emit 4p⁶4dⁿ \rightarrow 4dⁿ⁻¹4f +4p⁵4dⁿ⁺¹UTA, Sn¹⁰⁺–Sn¹³⁺ near 13.5 nm



Configuration Interaction effects are very important (Koike et al 2005 J. Elec. Spe Relat. Phen. 144, 1227, O'Sullivan and Faulkner 1994 Opt. Eng. 33, 3978)

> Optimum plasma conditions (Nishihara et al 2008)







Experimental Spectra of Xenon; Source Dependence



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Spectra from plasmas of elements with Z > 50

Spectra from elements with Z > 50 contain lines and an intense UTA due to $4p^{6}4d^{n} - 4p^{5}4d^{n+1}+4d^{n-1}4f$ ($0 \le n \le 9$) transitions. The effects of CI are to cause a transitions in successive ion stages to overlap in energy.

The degree of overlap improves with Z up to Z = 62.







Shorter wavelength sources

IOP PUBLISHING Phys. Scr. 80 (2009) 045303 (6pp)

EUV spectra of Gd and Tb ions excited in laser-produced and vacuum spark plasmas

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Interest in sources at 6.7 nm due to availability of Mo/B4C multilayer mirrors with a reflectivity of 40%

APPLIED PHYSICS LETTERS 97, 111503 (2010)

Rare-earth plasma extreme ultraviolet sources at 6.5-6.7 nm

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Figure 1. Spectra of gadolinium ions excited in the vacuum spark (upper trace) and in the laser-produced plasma (bottom trace).



Figure 2. Spectra of terbium ions excited in the vacuum spark (upper trace) and in the laser-produced plasma (bottom trace). *, 4f²-4f5d transition array in Tb XVIII classified in the present work.







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Most Important Stages?

IOP PUBLISHING

J. Phys. B; At. Mol. Opt. Phys. 43 (2010) 205004 (14pp)

OF PHYSICS B; ATOMIC, MOLECULAR AND OPTICAL PHYSICS doi:10.1088/0953-4075/43/20/205004

Tungsten spectra recorded at the LHD and comparison with calculations

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The most important transitions can be inferred from studies of W spectra.

They occur in Ag-like, Pd-like and Rh-like W^{27} +- W^{30+} . Sugar et al JOSA 10, 1321 (1993)

Gd XVIII-XX, Tb XIX - XXI

i.e. lons with 4d¹⁰4f, 4d¹⁰ and 4d⁹ ground states





Ag-like and Pd-like lines

Physica Scripta. Vol 26, 419-421, 1982

Resonance Lines in the Pd I Isoelectronic Sequence: **Tb XIX & XX** I VIII to Ho XXII J. Sugar and V. Kaufman Tb Laser: **Tb** XIX 3 J in 20 ns $\lambda = 1.06 \,\mu m$ Physica Scripta. Vol. 24, 742-746, 1981. Tb XIX $\Phi = (5-8) \times 10^{11} \, \text{Wcm}^{-2}$ Ag I Isoelectronic Sequence: Wavelengths and Energy Levels for Ce XII through Ho XXI and for W XXVIII Vacuum J. Sugar and V. Kaufman spark Tb XIX Tb XIX Laser Tb XX plasma (TTT 50 60 70 80 90 100 110 x.Å

Figure 2. Spectra of terbium ions excited in the vacuum spark (upper trace) and in the laser-produced plasma (bottom trace). *, 4f²-4f5d transition array in Tb XVIII classified in the present work.





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MARIE CURIE Ion populations and average ionization of a Gd plasma as a function of T_e computed with the CR model.





Power Density Requirements

The laser power density required lies in the range

 $2x10^{12} - 10^{13} \text{ Wcm}^{-2}$ @ $\lambda = 1.06 \text{ }\mu\text{m}$ $10^{11} - 10^{12} \text{ Wcm}^{-2}$ @ $\lambda = 10.6 \text{ }\mu\text{m}$

> Next Step: Calculations for emission from different ions



Fig. 5a: Approx. laser power density required for the RbI and PdI like isoelectronic sequences







Calculations for Gd and Tb

Singly excited configurations of Gd and Tb ions included in the calculations

Gd xvII - Gd xVIII	Gd XIX	Gd xx - Gd xxvIII	
Tb xviii - Tb xix	Tb xx	Tb xxi - Tb xxix	
$4p^{6}4d^{10}4f^{M}$	$4p^{6}4d^{10}$	$4p^{6}4d^{N}$	1
$4p^{6}4d^{9}4f^{M+1}$	$4p^{6}4d^{9}nl$	$4p^{6}4d^{N-1}nl$	
$4p^64d^94f^Mnl$		$4p^54d^{N+1}$	
$(n \leq 8, l \leq 3, M \leq 2)$	$(1 \le N \le 10)$		

Problem, low ion stages contain open 4f (and 5p) subshells, difficult to calculate.

(Kilbane and O'Sullivan JAP in press)



Energy levels of Gd¹⁶⁺ - Gd²⁷⁺ computed with the FAC code including CI. The number of states of each ion is shown at the top.





Effects of CI in Gd and Tb Spectra



Gd²⁰⁺ - Gd²⁵⁺, and Tb²¹⁺ - Tb²⁶⁺, spectra excluding CI (left) and including (right). 4d -4f (blue), 4p -4d (red) and all transitions (black)

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CI not as effective at spectral narrowing in highest stages







Line Strengths of Gd¹⁶⁺ - Gd²⁷⁺ and Tb¹⁷⁺ - Tb²⁸⁺ including Cl









Emission of Gd and Tb at various electron temperatures T_e



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Theoretical and experimental (Churilov et al) Gd and Tb spectra (a) (a) (b) Tb 115 4 (churilov et al) Gd and (churilov et al) Gd an



Other results will be discussed in:

Rare-earth Plasma Extreme Ultraviolet Sources at 6.5-6.7 nm for Next Generation Semiconductor Lithography (P6) Takeshi Higashiguchi, Utsunomiya University







$\Delta n = 0$, n = 4 UTA in other elements



Sfi

At low concentration get high spectral purity and low opacity







Ground configurations of ions in stages V – XVII for elements lanthanum through hafnium.

	V	VI	VII	VIII	IX	Х	XI	XII	XIII	XIV	XV	XVI	XVII
\mathbf{La}	$5s^{2}5p^{5}$	$5s^{2}5p^{4}$	$5s^25p^3$	$5s^{2}5p^{2}$	$5s^{2}5p$	$5s^{2}$	5s	$4d^{10}$	$4d^9$	$4d^8$	$4d^7$	$4d^6$	$4d^5$
Се	$5s^{2}5p^{6}$	$5s^{2}5p^{5}$	$5s^{2}5p^{4}$	$5s^{2}5p^{3}$	$5s^{2}5p^{2}$	$5s^25p$	$5s^2$	5s	$4d^{10}$	$4d^9$	$4d^8$	$4d^7$	$4d^6$
\mathbf{Pr}	$5p^{6}4f$	$5s^{2}5p^{6}$	$5s^{2}5p^{5}$	$5s^{2}5p^{4}$	$5s^{2}5p^{3}$	$5s^{2}5p^{2}$	$5s^24f$	$5s^{2}$	5s	$4d^{10}$	$4d^9$	$4d^8$	$4d^7$
Nd	$5p^{6}4f^{2}$	$5p^{G}4f$	$5s^25p^6$	$5s^{2}5p^{5}$	$5p^{3}4f$	$5p^24f$	$5s^{2}4f^{2}$	$5s^{2}4f$	$5s^{2}$	5 <i>s</i>	4d ¹⁰	$4d^9$	$4d^8$
\mathbf{Pm}	$5p^{6}4f^{3}$	$5p^{6}4f^{2}$	$5p^64f$	$5p^54f$	$5p^34f^2$	$5p^24f^2$	$5s^{2}4f^{3}$	$5s^{2}4f^{2}$	$5s^24f$	$5s^{2}$	4f	$4d^{10}$	$4d^9$
\mathbf{Sm}	$5p^{6}4f^{4}$	$5p^{6}4f^{3}$	$5p^{6}4f^{2}$	$5p^{5}4f^{2}$	$5p^{3}4f^{3}$	$5p^{2}4f^{3}$	$5s^24f^4$	$5s^{2}4f^{3}$	$5s^{2}4f^{2}$	$5s^24f$	5s4f	4f	$4d^{10}$
$\mathbf{E}\mathbf{u}$	$5p^{6}4f^{5}$	$5p^{6}4f^{4}$	$5p^{6}4f^{3}$	$5p^{5}4f^{3}$	$5p^{3}4f^{4}$	$5p^24f^4$	$5s^{2}4f^{5}$	$5s^{2}4f^{4}$	$5s^{2}4f^{3}$	$5s^{2}4f^{2}$	$5s4f^2$	$4f^2$	4f
Gd	$5p^{6}4f^{6}$	$5p^{6}4f^{5}$	$5p^64f^4$	$5p^{5}4f^{4}$	$5p^44f^4$	$5p^{2}4f^{5}$	$5s^{2}4f^{6}$	$5s^{2}4f^{5}$	$5s^{2}4f^{4}$	$5s^{2}4f^{3}$	$5s4f^3$	$4f^3$	$4f^2$
тъ	$5p^{6}4f^{7}$	$5p^{6}4f^{6}$	$5p^{6}4f^{5}$	$5p^{5}4f^{5}$	$5p^{4}4f^{5}$	$5p^{2}4f^{6}$	$5s^{2}4f^{7}$	$5s^{2}4f^{6}$	$5s^{2}4f^{5}$	$5s^{2}4f^{4}$	$5s4f^4$	$4f^4$	$4f^3$
$\mathbf{D}\mathbf{y}$	$5p^{6}4f^{8}$	$5p^{6}4f^{7}$	$5p^{6}4f^{6}$	$5p^{5}4f^{6}$	$5p^{4}4f^{6}$	$5p^{2}4f^{7}$	$5p4f^7$	$5s^{2}4f^{7}$	$5s^{2}4f^{6}$	$5s^24f^5$	$5s^24f^4$	$4f^5$	$4f^4$
$_{\rm Ho}$	$5p^{6}4f^{9}$	$5p^{6}4f^{8}$	$5p^{6}4f^{7}$	$5p^{5}4f^{7}$	$5p^44f^7$	$5p^{2}4f^{8}$	$5p4f^{8}$	$5s^{2}4f^{8}$	$5s^{2}4f^{7}$	$5s^{2}4f^{6}$	$5s^24f^5$	$4f^6$	$4f^5$
\mathbf{Er}	$5p^{6}4f^{10}$	$5p^{6}4f^{9}$	$5p^{6}4f^{8}$	$5p^54f^8$	$5p^44f^8$	$5p^24f^9$	$5p4f^{9}$	$5s^{2}4f^{9}$	$5s^{2}4f^{8}$	$5s^{2}4f^{7}$	$5s^24f^6$	$5s4f^{6}$	$4f^6$
Tm	$5p^{6}4f^{11}$	$5p^{6}4f^{10}$	$5p^{6}4f^{9}$	$5p^84f^8$	$5p^{4}4f^{9}$	$5p^24f^{10}$	$5p4f^{10}$	$5s^24f^{10}$	$5s^{2}4f^{9}$	$5s^{2}4f^{8}$	$5s^{2}4f^{7}$	$5s4f^7$	$4f^7$
Yb	$5p^{6}4f^{12}$	$5p^{6}4f^{11}$	$5p^{6}4f^{10}$	$5p^{6}4f^{9}$	$5p^{4}4f^{10}$	$5p^{3}4f^{10}$	$5p4f^{11}$	$5s^24f^{11}$	$5s^24f^{10}$	$5s^{2}4f^{9}$	$5s^{2}4f^{8}$	$5s4f^{8}$	$4f^8$
$\mathbf{L}\mathbf{u}$	$5p^{6}4f^{13}$	$5p^{6}4f^{12}$	$5p^{6}4f^{11}$	$5p^{6}4f^{10}$	$5p^{4}4f^{11}$	$5p^{3}4f^{11}$	$5p4f^{12}$	$5s^24f^{12}$	$5s^24f^{11}$	$5s^24f^{10}$	$5s^{2}4f^{9}$	$5s4f^{9}$	$4f^9$



Discrepancies between the current table and Carroll and O'Sullivan (1982) are highlighted in red.





Pd-like through Rblike spectra of La, Ce and Pr

Excluding CI (left) Including CI (right).

4d - 4f (blue), 4p - 4d (red) and all transitions (black)

(Kilbane and O'Sullivan PRA in press)







Pd-like through Rb-like spectra of Nd through Lu



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Maximum peak emission from $\Delta n = 0, 4 - 4$ UTAs in the lanthanides.

Dependence of UTA transition energies on Z. O'Sullivan and Carroll 1982(o), current work (x)









Conclusion

- $\Delta n = 0$, n = 4 4 UTA provide bright EUV sources in the 5 13.5 nm range
- •Broadly same physics as Sn sources, but ion stages with 4d subshell greater than half full now more important (opposite to Sn case)
- •Opacity an issue...low density targets or CO₂ plasmas
- •Strongest lines expected from Ag-like and Pd-like ions.
- •Emission bandwidth less than Sn and narrowest around Ce/Nd.









Thank You!

Acknowledgement: This work was supported by Science Foundation Ireland under Principal Investigator Grant 07/IN.1/I1771 and Marie Curie FP7 –People-IAPP Grant FIRE



