EUV emission from tin plasmas

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Why do we care about emission from tin? - lithography



A laser-produced Sn plasma emits strong radiation in a narrow band centered around 13.5 nm. This has high potential as an efficient EUV radiation source for use in the micro-electronics industry. The challenge is to make this efficient!



Why do we care about emission from tin? – atomic physics

- We are always interested in measurements that will help us validate our atomic physics models
- The plasma conditions (around 30 eV, 0.1% or less of solid density) lead to emission from Sn ions that have between 7-14 electrons removed
- The challenge for theory is to accurately describe these due to strong configuration-interaction effects in the atomic structure involving 4p-4d and 4d-4f transitions
- AND construct a plasma model that can efficiently predict the ionization balance & emission from such a plasma







The LANL suite of atomic modeling codes

CATS: Cowan Code

RATS: relativistic

ACE: e⁻ excitation

GIPPER: ionization

http://aphysics2.lanl.gov/tempweb

fine-structure config-average UTAS **MUTAs** energy levels gf-values e⁻ excitation e⁻ ionization photoionization autoionization

LTE or NLTE populations

spectral modeling emission absorption transmission power loss

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Atomic structure calculations - additions

- Our version of **Cowan's code** (**CATS**) has been extensively re-written into Fortran90 and parallelized so that each Jpi symmetry is on a separate processor of a parallel machine. This speeds up runtime considerably
 - Also, when computing dipole matrix elements (for gf values), each J/J' combination is also placed on a separate processor.
 - The memory requirements and runtime requirements are still considerable sometimes 100s GB RAM memory and runtimes approaching one week per ion stage
 - We have a dedicated workstation to perform atomic structure calculations, with a large hard drive
- We also include a **2-mode** option. This modification allows a user-specified number of configurations to be treated with full configuration-interaction (CI), while any other configurations are treated through intermediate-coupling (IC).
 - IC is much cheaper, computationally, and many (up to 10⁴ configurations or more) may be treated this way. This provides enough excited configurations to ensure a well-converged partition function when computing an opacity.
- The **scale factors** used in CATS are defaulted to an option that was designed to scale with Z and ion stage. We have used these for the ground state to excited state calculations, but modified the scale factors for the excited-state to excited state calculations.



Scale factors? According to Cowan, these are necessary to improve agreement with experimentally measured transition wavelengths – and are used to account for the *"infinity of small perturbations"* that are necessarily omitted in practical calculations





Defining the atomic physics problem – accuracy and quantity of data both issues



Our FS model contains all the configurations we expect to be important for CI effects:

- Sn 14+: 114 cfgs; 94115 levels
- Sn 13+: 135 cfgs; 273330 levels
- Sn 12+: 94 cfgs; 355742 levels
- Sn 11+: 81 cfgs; 259181 levels

Just these 4 ionization stages generate ~ 30 billion dipole-allowed transitions!







 How does the position and magnitude of the absorption features change as more configurations are added?



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Model 1: 3 cfgs; 141 levels; 547 transitions Model 2: 7 cfgs; 1696 levels; 282216 transitions Model 3: 12 cfgs; 40317 levels; 125M transitions Model 4: 18 cfgs; 48687 levels; 184M transitions Model 5: 21 cfgs; 85733 levels; 595M transitions Model 4a: 33 cfgs; 50561 levels; 200M transitions

Addition of excitations to n=5 also modify the main feature





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Opacity (cm^2/g) .5e+06 5e+05 le+06 Model 1: 3 cfgs; 141 levels; 547 transitions Model 1: Model 2: 7 cfgs; 1696 levels; 282216 transitions Model 3: 12 cfgs; 40317 levels; 125M transitions Model 4: 18 cfgs; 48687 levels; 184M transitions Model 5: 21 cfgs; 85733 levels; 595M transitions Model 4a: 33 cfgs; 50561 levels; 200M transitions Model 4b: 53 cfgs; 138499 levels; 1593M transitions Wavelength (nm Addition of excitations to n=5 also modify the main prominent feature 5 Model Model Mode Model Model Model Model 16

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However: closer examination of the individual features shows that absolute convergence is difficult to obtain!







Wavelength

 How does the position and magnitude of the absorption features change as more configurations are added?

Model 1: 3 cfgs; 141 levels; 547 transitions Model 2: 7 cfgs; 1696 levels; 282216 transitions Model 3: 12 cfgs; 40317 levels; 125M transitions Model 4: 18 cfgs; 48687 levels; 184M transitions Model 5: 21 cfgs; 85733 levels; 595M transitions Model 4a: 33 cfgs; 50561 levels; 200M transitions Model 4b: 53 cfgs; 138499 levels; 1593M transitions FSCI-n5e: 94 cfgs; 355742 levels; 10B transitions

- FSCI model appears reasonably well converged with respect to main absorption feature
- This is then repeated for all other relevant ion stages
- Calculations are extended using our "2-mode" method to include contributions from other, higherying, transitions

Opacity (cm^2/g) .5e+06 2e+06 5e+05 le+06 Model 1:4p FSCI-n5e: Model 4b Model 4b: Mode 12 Model Mode Mode Mode Model Model Aodel 16

Conclusions & Future Work

- Emission spectra of Sn plasma at moderate temperatures is very demanding to compute
- Configuration-interaction is very important in such species, making the structure calculations complex and demanding
- Multiply excited states are found to make significant contributions to the plasma, even at moderate densities
- Agreement with laser-produced plasma measurements is very encouraging
 - Comparisons also show that taking into account radiation transport effects is important
- We continue towards our ultimate goal of a predictive set of opacity and emissivity calculations for such



